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A Study of Muon Pair Production and Evidence for Tau Pair Production in Photon-Photon Collisions at LEP

The OPAL Collaboration

Abstract

We have studied the reactions $\gamma\gamma \rightarrow \mu^+\mu^-$ and $\gamma\gamma \rightarrow \tau^+\tau^-$ by measuring the properties of events of the types $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ and $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ at e^+e^- centre-of-mass energies between 88 and 94 GeV. The data sample corresponds to an integrated e^+e^- luminosity of 40.4 pb^{-1} collected by the OPAL detector at LEP. The QED structure function F_2 is extracted from a sample of 1462 $\gamma\gamma \rightarrow \mu^+\mu^-$ events in which one photon is off the mass shell (single-tagged). The Q^2 range for these massive photons is $4 - 400 \text{ GeV}^2$, with an average Q^2 of 8 GeV^2 . We have observed 48 $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ events in which both final state electrons are detected. In the sample of single-tag events we have identified 34.9 ± 6.7 events as due to the reaction $\gamma\gamma \rightarrow \tau^+\tau^-$, on an estimated background of 5.1 ± 2.7 events. In all cases the measured event distributions agree with QED calculations.

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1 Introduction

Two-photon reactions have been studied at e^+e^- colliders for some time [1,2] and measurements of the reaction $\gamma\gamma \rightarrow \mu^+\mu^-$ have been made at e^+e^- centre-of-mass energies of up to 61.4 GeV [3-10]. The two-photon reactions $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ and $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ are pure $\mathcal{O}(\alpha^4)$ QED reactions and are, therefore, in principle well understood.

In this paper, we present a measurement of the reaction $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ under two conditions: *i*) single tag, i.e. one of the final state electrons was observed, and *ii*) double tag, i.e. both final state electrons were observed. In both cases, we required the two final state muons to be observed. Untagged events, those in which both final state electrons escaped detection close to the beam axis were not considered.

This paper also describes the first direct observation of tau pair production in photon-photon collisions. The high energy of the LEP beams makes the $\gamma\gamma$ mass range needed for $\tau^+\tau^-$ production more accessible. This feature of LEP and the excellent particle identification properties of the OPAL detector enable us to measure the reaction $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ by selecting final states with one electron and one muon or charged hadron, in addition to the tag. In order to reduce backgrounds only single-tag events were considered in this study.

The data sample consists of the data collected from 1990 to 1992, which amounts to 40.4 pb^{-1} of integrated e^+e^- luminosity in the centre-of-mass energy range between 88 and 94 GeV.

2 Theory and Monte Carlo Simulations

The four main diagrams contributing at the lowest order to the reactions $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ and $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ are shown in Figure 1. The classical “two-photon” diagram is the one shown as “multiperipheral”, Figure 1a. We compare the data to the following Monte Carlo simulations:

1. The single-tag $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ events were simulated using the four-fermion generator written by Berends, Daverveldt and Kleiss [11] (BDK), which includes all four diagrams. We have used the version modified by Daverveldt to include the contributions coming from the annihilation through the Z^0 . It is optimised for the generation of untagged events and usable for single-tag events, but it is very inefficient in generating double-tag events.
2. For the simulation of double-tag data, we used the FERMISV [12] generator, which includes the same diagrams, and also initial and final state radiation.
3. The $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ events were simulated using the Vermaseren generator [13] which includes the diagrams of Figures 1a and 1b. The decay of the tau leptons was simulated with the routines from the JETSET library [14,15].

The events from all three generators were processed by the OPAL detector simulation code [16] and the standard offline reconstruction code. All events were generated at a centre-of-mass

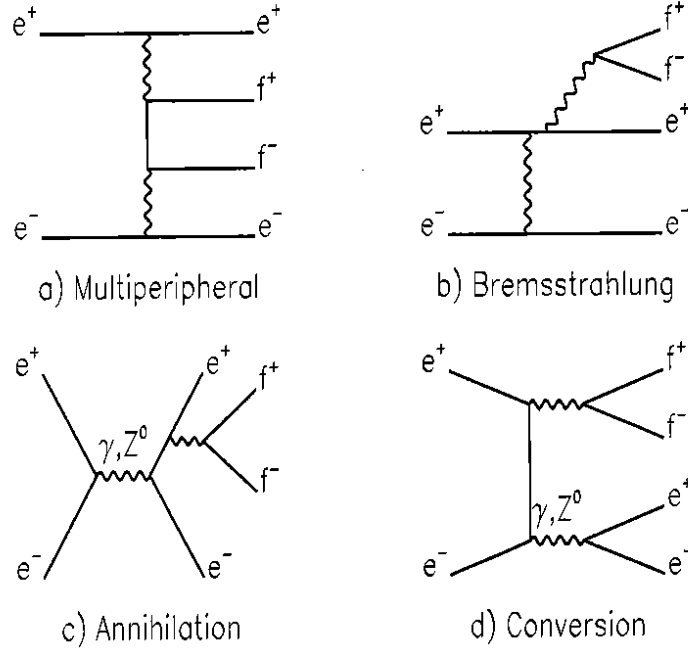


Figure 1: The four main diagrams contributing at the lowest order to the processes $e^+e^- \rightarrow e^+e^-f^+f^-$, with $f = \mu, \tau$. In b) both initial and final state radiation are included.

energy corresponding to the Z^0 mass, where 33.1 of the 40.4 pb^{-1} of data were taken. The BDK generator was used to verify that the bremsstrahlung diagram contributes 2% of the events to the single-tag $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ channel. The annihilation and conversion diagrams both contribute less than 0.1%, and were neglected in the single-tag analysis.

For comparing the photon structure function F_2 obtained from the data with the QED prediction, we have used TWOGEN, a Monte Carlo which generates only the two-photon flux factor N which can then be folded with any structure function [17,18]. The differential cross section for single-tag events can be written as:

$$\frac{d^3\sigma}{dx dQ^2 dz} = \frac{4\pi\alpha^2}{Q^4 x} \{ (1-y)F_2(x, Q^2) + xy^2 F_1(x, Q^2) \} \frac{dN}{dz}. \quad (1)$$

Here Q^2 is the momentum transfer squared of the photon, tagged by the measurement of the scattered electron. The kinematic variable x is defined as $x = Q^2/(Q^2 + W^2)$, where W is the invariant mass of the muon pair, $y \approx 1 - E_{\text{tag}}/E_{\text{beam}}$ and $z = 1 - E_{\text{untag}}/E_{\text{beam}}$. The energies of the tagged and the untagged electron are E_{tag} and E_{untag} , respectively. The untagged electron is assumed to have $Q^2 = 0$. The QED structure functions of the photon are predicted to be [19]:

$$F_1(x, Q^2) = \frac{\alpha}{2\pi} \left\{ [x^2 + (1-x)^2 + 4m_\mu^2 \frac{W^2 - 2m_\mu^2}{(W^2 + Q^2)^2}] \log \left[\frac{W}{2m_\mu} + \left(\frac{W^2}{4m_\mu^2} - 1 \right)^{1/2} \right]^2 - \left[(1-2x)^2 + \frac{4m_\mu^2 W^2}{(W^2 + Q^2)^2} \right] \left(1 - \frac{4m_\mu^2}{W^2} \right)^{1/2} \right\}, \quad (2)$$

and

$$F_2(x, Q^2) = 2xF_1(x, Q^2) + \frac{4\alpha}{\pi} x^2 \left\{ \left(1 - x \right) \left(1 - \frac{4m_\mu^2}{W^2} \right)^{1/2} - \frac{2m_\mu^2}{W^2 + Q^2} \log \left[\frac{W}{2m_\mu} + \left(\frac{W^2}{4m_\mu^2} - 1 \right)^{1/2} \right]^2 \right\}. \quad (3)$$

The contribution from the structure function $F_1(x, Q^2)$ is negligible for $y < 0.25$. In this paper, we shall only be concerned with the determination of $F_2(x, Q^2)$.

3 The OPAL Detector

As the OPAL detector has already been described in detail [20], we restrict ourselves here to an overview of the main components. The coordinate system of the detector is defined with the z -axis along the electron beam direction, ϑ and φ being the polar and azimuthal angles. Tracking of charged particles is mainly performed by the jet chamber, a large volume drift chamber, 4 m long and 3.7 m in diameter, divided into 24 azimuthal sectors each containing a plane of 159 sense wires. The charge measured by the sense wires is used to obtain dE/dx , the specific energy loss of the particle in the detector gas. Together with the momentum, this gives a measurement of the mass of the particle, leading to an e - π separating power of more than two standard deviations for momenta up to 12 GeV [21]. The jet chamber, the vertex detector, and the chambers measuring the z coordinates of particles as they leave the jet chamber, are positioned inside a solenoidal coil which provides an axial magnetic field of 0.435 T. It is surrounded by a time-of-flight counter array (TOF), a presampler and an electromagnetic calorimeter.

The electromagnetic calorimeter consists of a barrel and two endcap arrays of lead glass blocks. Each block subtends a solid angle of approximately 40×40 mrad², with a thickness of over 24 radiation lengths (X_0) in the barrel region and typically $22X_0$ in the endcaps. The energy resolution for electrons of 45 GeV is typically $\sigma_E/E \approx 3\%$. Outside the electromagnetic calorimeter is the instrumented return yoke of the magnet, forming the hadron calorimeter, and beyond this are four layers of the outer muon detectors. The hadron calorimeter, which consists of nine layers (eight layers in the endcap) of streamer tubes interleaved with the iron slabs of the magnet return yoke, is read out via 4 mm wide strips spaced 1 cm apart, and projective towers formed by pads. These strips provide tracking in the (r, φ) plane, while the towers measure the energy deposited in coarse (φ, z) bins. The material in front of the hadron calorimeter corresponds to about two hadronic interaction lengths. The hadron calorimeter itself accounts for about 4.8 interaction lengths. A muon with a transverse momentum $p_t = 1$ GeV will on the average traverse six planes of the hadron calorimeter.

The forward detectors consist of several sub-detectors of which three are used in this analysis: a lead-scintillator sandwich calorimeter, a set of proportional tube chambers embedded in the calorimeter, and two planes of drift chambers upstream of the calorimeter.

The forward detector calorimeter measures the energy deposited by a photon or an electron with a resolution of $\Delta E/E = 18\%/\sqrt{E}$, E in GeV. Careful calibration of the energy scale keeps the fluctuations in the measured shower energy between LEP fills to less than 0.5%. The calorimeter also measures the shower position by virtue of its azimuthal division into 16 segments and separate inner and outer edge readouts. The polar angle resolution varies between 1 and 10 mrad, being best near the inner edge of the calorimeter. The resolution in azimuth varies between 3.5 and 35 mrad, being best at the segment boundaries. Furthermore, the calorimeter gives shower development information through longitudinal segmentation into presampler ($4X_0$) and main ($19X_0$) sections. The acceptance of the calorimeter extends from 39 to 155 mrad, and is essentially complete in azimuth. The tube chambers are situated after the presampler section and measure the shower position in ϑ and φ . The resolution for 45 GeV electrons is 1.3 mrad in ϑ and 19 mrad in φ . The tube chamber acceptance extends from 50 to 135 mrad in ϑ and covers 95% in azimuth. The precisely located drift chambers are used to determine the absolute radial position of the tube chambers to a precision of 135 μ m. The distance between the forward detectors at each side of OPAL is known to a precision of 0.5 mm.

4 The Event Selection

In the events under study we distinguish: *i*) the high-energy electrons (tags), and *ii*) the low-energy muons or the decay products of the tau pair. These can be handled independently.

The pattern recognition program reconstructed charged tracks from strings of at least eight hits in the jet chamber ($\vartheta > 200$ mrad). A good track was defined in this analysis as having: at least 30 hits registered in the jet chamber, the first one at a radius smaller than 75 cm; a distance of closest approach to the nominal beam axis of less than 1 cm, and $|z| < 20$ cm at this point. The momentum of the track had to be between 300 MeV and 20 GeV with a component transverse to the beam axis (p_t) greater than 100 MeV.

If a track failed one or more of these criteria, it could be defined as being of medium quality if the radius of its first measured point was at a distance of less than 50 cm from the beam axis and it had at least 30 jet chamber hits associated with it. As a result, medium quality tracks would include low energy photon conversion tracks and badly measured low angle tracks, but not tracks which scattered back from the pressure vessel or magnet coil. The latter, and other tracks not classified as of medium or good quality were ignored.

An electromagnetic calorimeter cluster was defined as an energy deposition of more than 100 MeV in the barrel lead glass array, or more than 250 MeV in the endcap.

In the first stage of the selection the presence of a tag was established. A tag in the forward detector is defined as any electromagnetic energy deposition E_{tag} exceeding half the beam energy E_{beam} . We expect most tags to be observed in the forward detector because the tag angle spectrum is strongly peaked at small angles. In order to have well measured energy values, we required the polar angle ϑ_{tag} of this tag to be between 47 and 120 mrad. No distinction between photons and electrons was made in the forward detector. In order to ensure that the untagged electron in single-tag events had $Q^2 \approx 0$, we rejected events with additional energy depositions of less than $E_{\text{beam}}/2$ in the forward detector.

A tag with a polar angle above 200 mrad was defined using electromagnetic clusters in the barrel and endcap lead glass arrays. Such a cluster was called a tag if its energy exceeded half the beam energy and at least one charged track of any quality was found in a cone of 15° half-opening angle centred on the cluster. Additional charged tracks within this cone were ignored. No requirement was made on the charge of the tag.

In the next step of the selection only tracks not assigned to tags were considered for event selection. The events were required to have two and only two good tracks with opposite charges. Events with additional medium quality tracks were rejected, as were events with reconstructed photon conversions. At least one of the tracks had to have a momentum greater than 1 GeV and a polar angle ϑ satisfying $|\cos \vartheta| < 0.72$. The other charged track had to have a momentum larger than 300 MeV and a polar angle satisfying $|\cos \vartheta| < 0.95$.

The reaction $e^+e^- \rightarrow \tau^+\tau^-$ was found to be a background to events in which the tag was measured at $\vartheta > 200$ mrad. This background was rejected by the requirement that $\vartheta_{\text{untag}} < 200$ mrad. Here ϑ_{untag} is the polar angle of the untagged electron, reconstructed from the momentum vectors of the tagged electron and the two charged tracks. The value of 200 mrad allows for uncertainties in the momentum measurements.

All events were required to have initiated the detector read-out with a trigger consisting of a coincidence between an energy deposition of more than 15 GeV in one of the two forward detector modules and a charged track in the barrel region of the jet chamber. The efficiency of this trigger for events satisfying all cuts was determined by means of other, independent, triggers to be 95.3 ± 0.6 %.

5 The Event Identification

Identification of the final state particles was required for further selection. Identification criteria were defined for electrons and muons. Figures 2a and b show the overall efficiencies for identifying electrons and muons as functions of their momentum. The slight disagreement between data and Monte Carlo at high electron momenta has no influence on the results.

The electron identification was based on two measured quantities: the ratio of electromagnetic energy measured in the lead-glass calorimeter to the momentum measured in the central detector E/p , and the energy loss dE/dx in the jet chamber. The E/p selection was used for momenta greater than 3 GeV: a charged particle was called an electron if the electromagnetic cluster associated with it had an energy E satisfying $E/p > 0.8$. Between 0.3 and 3.0 GeV a particle was identified as an electron if $9.0 < dE/dx < 12.0$ keV/cm. For both momentum ranges, it was required that no signal above 200 MeV was observed in the region of the hadron calorimeter associated with the charged track.

Charged particles were identified as muons only if their momentum exceeded 1 GeV. The electromagnetic energy associated to the charged track was required to be less than 1.5 GeV. A signal was required in the muon chambers or in the hadron calorimeter strips. In the muon chambers, a track segment with at least three hits had to be reconstructed. In the hadron calorimeter, at least three layers had to show a strip signal associated with the track, of which one had to be at least in the fourth layer. The ratio of the total number of strips to the number of layers hit had to be less than three. This last requirement was to reject hadrons which penetrated several layers before initiating a shower. In order to satisfy the criteria for muon identification, a muon impinging perpendicularly on the hadron calorimeter had to traverse the lead glass array and at least three layers of iron, amounting to a total of four hadronic interaction lengths.

6 Single-Tag Mupair Events

We identified the $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ final state by requiring the presence of one identified muon and rejecting events with identified electrons other than the tag. With these cuts we selected 1462 events, of which 48 events had a tag with $\vartheta_{\text{tag}} > 200$ mrad.

6.1 Backgrounds and Systematic Errors

The process $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ constitutes the dominant background to $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$. Using the Vermaseren Monte Carlo we have estimated the background for all τ decay channels to be 24.0 ± 2.4 events. Cross sections for continuum production and resonance formation in $\gamma\gamma \rightarrow \pi^+\pi^-$ are much lower than the $\gamma\gamma \rightarrow \mu^+\mu^-$ cross section and the pions are not expected to pass the muon selection criteria described above. The reaction $e^+e^- \rightarrow \tau^+\tau^-$ contributes 0.2 ± 0.2 event to the events with $\vartheta_{\text{tag}} > 200$ mrad.

The systematic errors in the reconstruction and selection efficiency were derived from the comparison between the data and the BDK Monte Carlo. We varied the cuts and compared the ratios of the numbers of events in the data and the Monte Carlo. We have estimated the errors due to the cuts in three intervals of x . Adding the errors in quadrature, we obtained systematic errors for the efficiency of 8.2% for $x < 0.4$, 11.9% for $0.4 < x < 0.7$ and 8.2% for $x > 0.7$. An error of 1% for the uncertainty in the luminosity measurement [22] was also included.

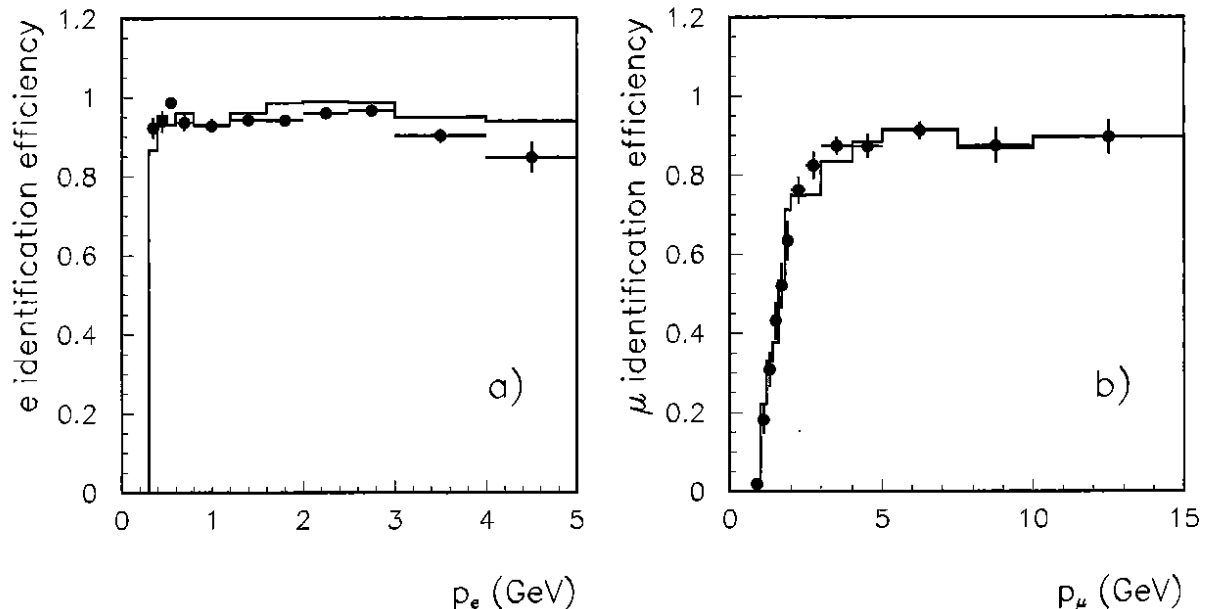


Figure 2: a) The efficiency for identifying an electron in a single-tag event in which the second particle was positively identified as an electron. b) The efficiency for identifying a muon in a single-tag event in which the second particle was positively identified as a muon. The efficiencies are plotted as functions of the particle momenta. The points represent the data, the solid lines represent the Monte Carlo.

We have studied the effect of initial state radiation by adding to the TWOGEN Monte Carlo the possibility of generating one bremsstrahlung photon collinear with the beam axis. We found that this changed the acceptance with our cuts by less than 1% and we have therefore neglected the effect of initial state radiation in this paper.

Taking into account the systematic errors and the statistical error of the Monte Carlo sample, we expect 1418 ± 142 events from the BDK Monte Carlo, of which 46 ± 7.4 with a tag at $\vartheta > 200$ mrad. The average values of x and Q^2 of the high- Q^2 events are 0.65 (RMS 0.64) and 127.8 (RMS 125.9), respectively, to be compared with 0.68 (RMS 0.67) and 140.4 (RMS 138.0) in the BDK Monte Carlo.

6.2 The QED Structure Function

The QED structure function F_2 was extracted from the data using the following procedure:

1. Single-tag $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ events were generated with the BDK Monte Carlo.
2. The generated events were subjected to the detector simulation and to the analysis cuts, and were weighted by a factor $\alpha/F_2(x, Q^2)$.
3. The x -distribution of the data was then divided by the Monte Carlo distribution of the weighted events. By doing this, we divided out the detector effects and also the photon flux factor, and obtained a distribution for F_2 , which can be compared directly with QED calculations.

The result is shown in Figure 3. The solid points are the measured values of F_2 as a function of x for the tagging range $\vartheta_{\text{tag}} > 47$ mrad and for $y < 0.25$. This corresponds to an

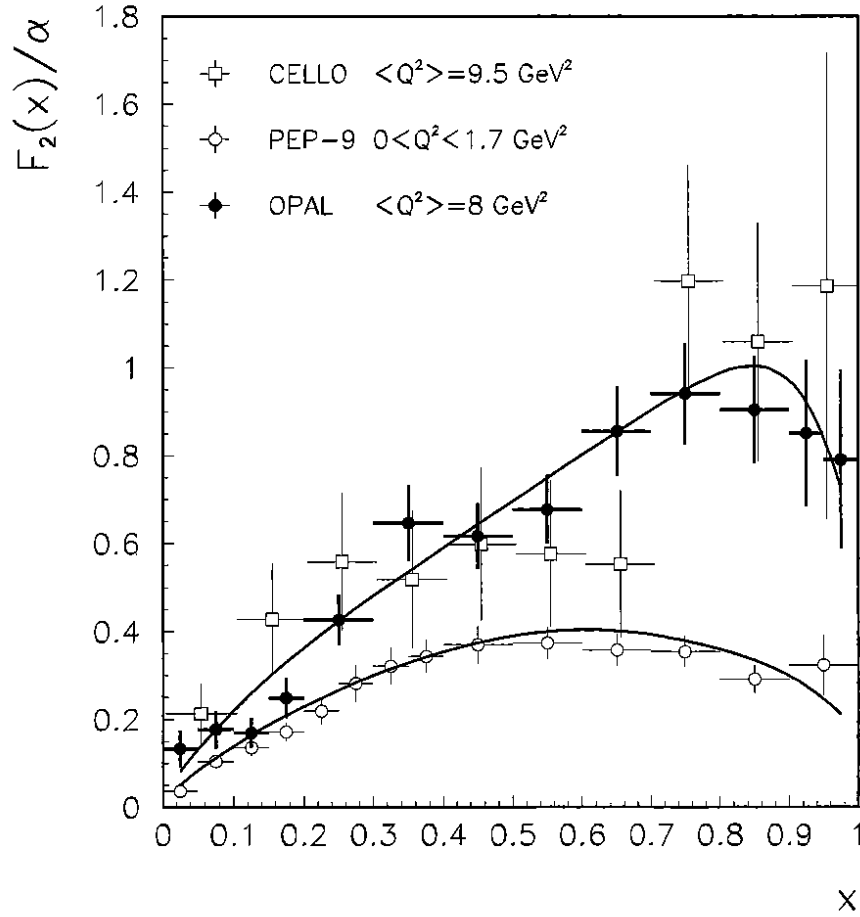


Figure 3: The measured values of the structure functions F_2 as functions of x for our measurement (full circles), the PEP-9 measurement (open circles) and the CELLO measurement (open squares). The solid lines are the QED expectations; the expectations for the CELLO and OPAL data are identical within the systematic errors.

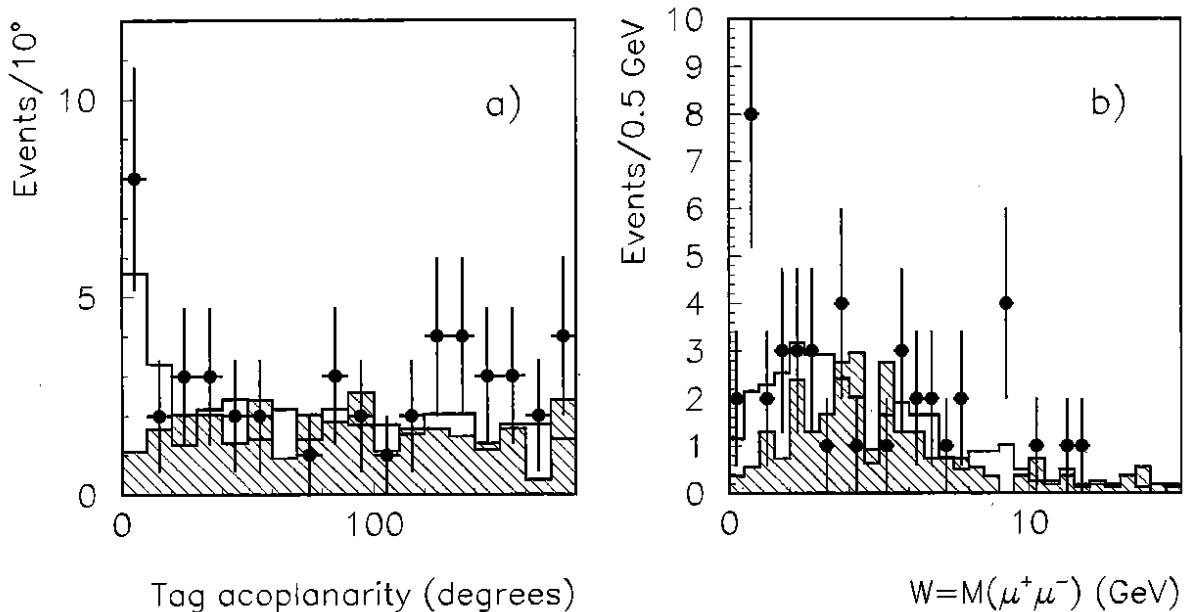


Figure 4: a) The distribution of double-tag $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ events as a function of the acoplanarity between the two final state electrons. The points represent the accepted events, the lines the Monte Carlo simulations (open: FERMISV, shaded: Vermaseren), both normalised to the integrated luminosity. b) The distribution of events as a function of W , the effective muon-muon mass. The data points and the histograms have the same meaning as in a). All errors are statistical only.

average Q^2 (calculated from the generated events) of 8.0 GeV^2 . The open circles are the PEP-9 measurements [4] for $0 < Q^2 < 1.7 \text{ GeV}^2$. The open squares are the CELLO measurements [3] at $< Q^2 >= 9.5 \text{ GeV}^2$. The solid lines represent the QED expectations. They were obtained by generating events according to the QED expression for F_2 with a two-photon luminosity generator as described in [17,18]. The expectation for the CELLO data is almost identical to the one shown for our data. The errors shown are statistical only. Our measurement is corrected for the estimated background from $\gamma\gamma \rightarrow \tau^+\tau^-$ events. It is in good agreement with the expectation.

7 Double-Tag Mupair Events

We found 48 events in which both final state electrons were detected, either in the forward or in the central detector. We expect 36.9 ± 5.0 such events from the FERMISV Monte Carlo for the same integrated luminosity. In the kinematic region of the double-tag events the contributions of the annihilation and conversion diagrams become significant (9.3% and 2.8%, respectively). This can be seen in the distribution of the events as a function of the acoplanarity of the tags, defined as $\pi - \varphi_{\text{tag1}} - \varphi_{\text{tag2}}$, shown in Figure 4a. The distribution is uniform, except for very small acollinearities which have an excess of events. This behaviour is reproduced by the FERMISV Monte Carlo (open histogram), but not by the Vermaseren Monte Carlo, which does not include the annihilation and conversion diagrams. Figure 4b shows the distribution of events as a function of the invariant $\mu^+\mu^-$ mass. We see a slightly better agreement with the data in the FERMISV Monte Carlo than in the Vermaseren Monte Carlo, especially at the low masses. We attribute the slight excess of events in the mass bin between 0.5 and 1 GeV to

a statistical fluctuation. We do not observe the excess of events around $W = 300$ MeV which was reported by the AMY collaboration [10].

As in the single-tag case, the reaction $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ is expected to be the dominant background to the double-tag $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ events. Using the Vermaseren Monte Carlo we have estimated this background to be 1.7 ± 0.6 events. No Monte Carlo events of the reaction $e^+e^- \rightarrow \tau^+\tau^-$ passed our cuts.

8 The Observation of $\gamma\gamma \rightarrow \tau^+\tau^-$ Events

The selection criteria for identification of the $\tau^+\tau^-$ final state were designed for events in which one tau decayed into an electron and the other tau into a muon or into a hadron with possibly additional neutral particles. This very characteristic decay mode represents about 25% of the total $\tau^+\tau^-$ decays. For any other tau decay mode, the backgrounds from other two-photon reactions were too large for our cuts.

The final selection criterion is discernable in Figure 5 which shows the distribution of tracks as a function of the particle identification variable T_R in the case that the other track in the event was identified as an electron. Here, $T_R \equiv [T_m - T_\mu(p)]/[T_e(p) - T_\mu(p)]$, T_m is the measured energy loss rate in the jet chamber (dE/dx), and $T_e(p)$ and $T_\mu(p)$ are the expected energy loss rates for an electron and a muon with momentum p , respectively. The points with error bars are the measured values; the shaded histogram corresponds to a Monte Carlo simulation of the $\gamma\gamma \rightarrow \tau^+\tau^-$ process with the Vermaseren Monte Carlo, normalised to the integrated luminosity of the data sample and corrected for the triggering efficiency. The estimated contribution from the background was added to the Monte Carlo. The smooth lines represent fits of Gaussian distributions to the data. The resolution obtained from the fit is 3.5%, in good agreement with the expectation [21]. The peak at about $T_R = 0$ corresponds to the $\gamma\gamma \rightarrow \tau^+\tau^-$ events accepted by our selection. The peak at $T_R = 1$ in the Monte Carlo distribution is due to events in which both tau leptons decay to electrons. The cut was applied at $T_R = 0.3$, which yields 43 events in the data.

8.1 Backgrounds for $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ Events

We have considered several sources of backgrounds for the $e^+e^- \rightarrow e^+e^-\tau^+\tau^-$ events:

1. The background due to the tail in the T_R distribution of $\gamma\gamma \rightarrow e^+e^-$ events was estimated by counting the number of tracks with $T_R > 1.7$ in Figure 5. We found 2 events and expect statistically that the contamination of events with $T_R < 0.3$ is the same. We found no $\gamma\gamma \rightarrow e^+e^-$ Monte Carlo events satisfying all cuts.
2. The background due to the reaction $\gamma\gamma \rightarrow \mu^+\mu^-$ was estimated from a Vermaseren Monte Carlo sample of $\gamma\gamma \rightarrow \mu^+\mu^-$ events to be 1.9 ± 1.1 events.
3. The backgrounds from events with larger multiplicities or with photon conversions in which one or more tracks were not reconstructed, was estimated by counting the number of events satisfying all cuts except the charge balance. We found one such event. This estimate includes the contribution from beam-gas interactions. In addition, the absence of beam-gas events is confirmed by the fact that the origin of the charged tracks (measured along the beam axis) corresponds within the experimental errors to the interaction point. In beam-gas events this origin is uniformly distributed along the beam axis.

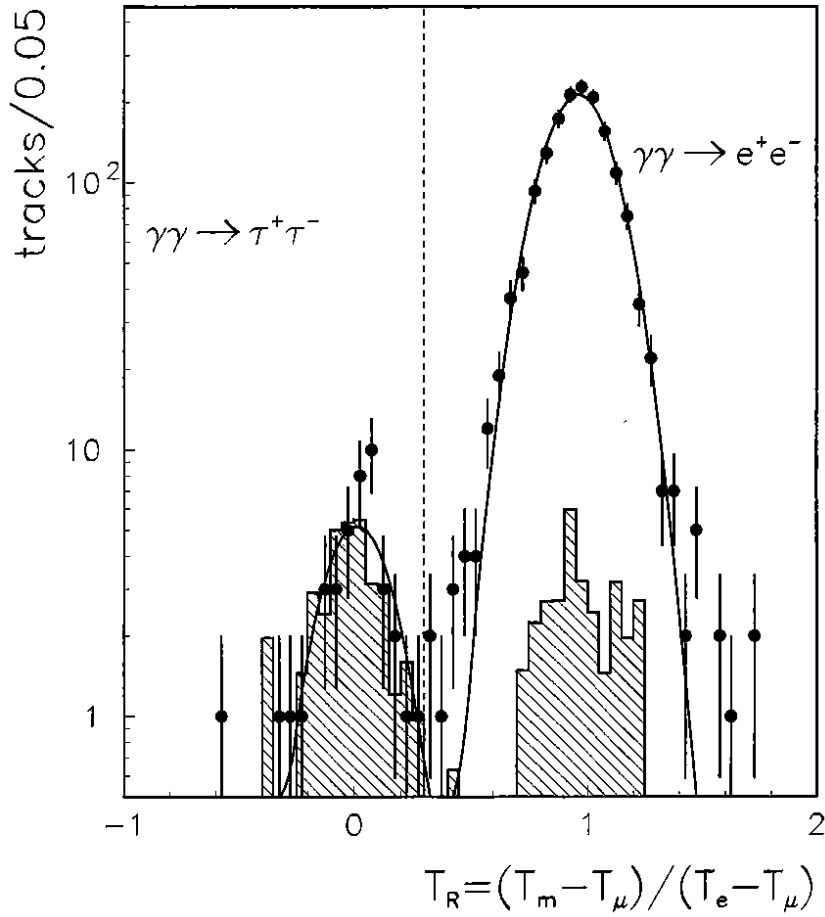


Figure 5: The distribution of tracks as a function of T_R in events in which the other track is identified as an electron. The points with error bars are the measured data, the shaded histogram corresponds to a Monte Carlo simulation of the $\gamma\gamma \rightarrow \tau^+\tau^-$ process with the Vermaseren Monte Carlo, normalised to the integrated luminosity of the data sample. The estimated background is added to the Monte Carlo histogram. The smooth lines are fits of Gaussian distributions to the data. The dashed line indicates where the cut was applied.

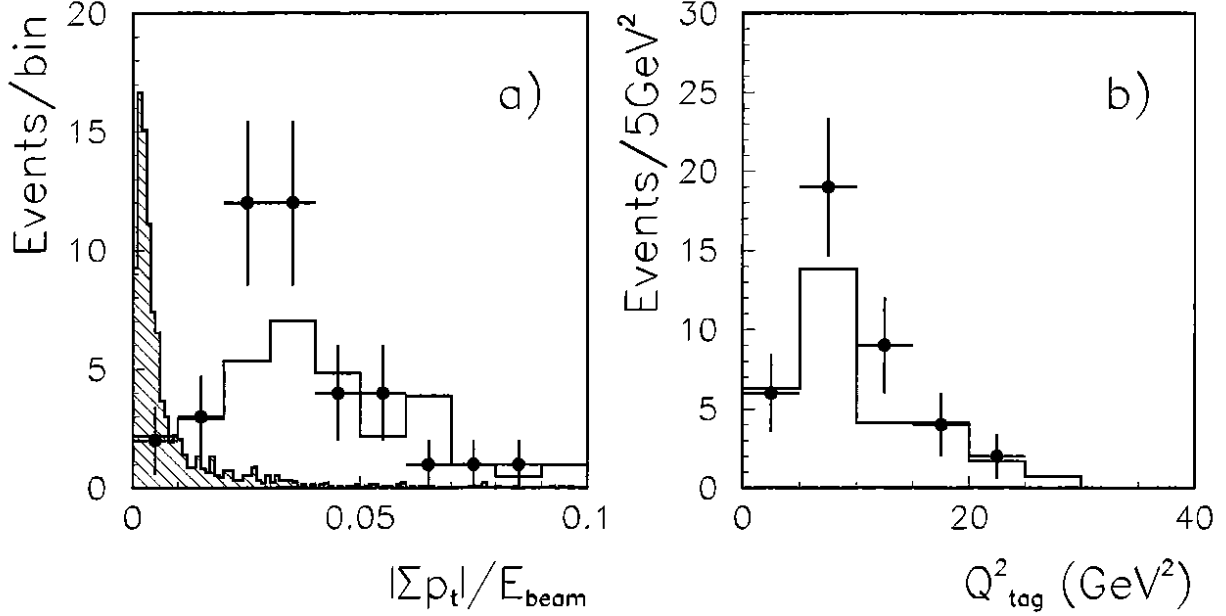


Figure 6: a) The distribution of $\gamma\gamma \rightarrow \tau^+\tau^-$ events as a function of the scaled transverse momentum sum, including the tagged electron. The points represent the accepted $\gamma\gamma \rightarrow \tau^+\tau^-$ events, the open histogram the $\gamma\gamma \rightarrow \tau^+\tau^-$ Monte Carlo, normalised to the integrated luminosity. The shaded histogram shows the same quantity for $\gamma\gamma \rightarrow \mu^+\mu^-$ events, with arbitrary normalisation. b) The distribution of $\gamma\gamma \rightarrow \tau^+\tau^-$ events as a function of Q^2 for the accepted events in the data (points with error bars) and in the Monte Carlo (solid line). All errors are statistical only.

4. The $e^+e^- \rightarrow \tau^+\tau^-(\gamma)$ Monte Carlo contributes 0.2 ± 0.2 event to the $\gamma\gamma \rightarrow \tau^+\tau^-$ sample.
5. The reactions $\gamma\gamma \rightarrow K^+K^-$ and $\gamma\gamma \rightarrow p\bar{p}$ can contribute to the sample if one of the kaons or protons lies in the dE/dx overlap regions e-K or e-p, and the other is identified as a muon or charged hadron. In two events, both particles were compatible with being kaons. The particles originally identified as electrons in these events had valid TOF signals which confirmed the kaon hypothesis. In the same way, one event was found to be compatible with $\gamma\gamma \rightarrow p\bar{p}$. We removed the events from the sample, and added the statistical error of ± 1.7 to the error in the background.

Adding these numbers, we estimate a background of 5.1 ± 2.7 events in the $\gamma\gamma \rightarrow \tau^+\tau^-$ event sample.

With the cuts described above we have selected 40 events in the data, with an estimated contamination of 5.1 ± 2.7 events due to various background reactions. We assume that the systematic error is of a similar magnitude as the systematic error established for the $\gamma\gamma \rightarrow \mu^+\mu^-$ acceptance. Including again a systematic error of 1% on the luminosity measurement for this sample we expect 32.7 ± 4.8 events from the Vermaseren Monte Carlo with the same cuts as applied to the data. The error also includes the statistical error of the Monte Carlo sample.

In Figure 6a the distribution of the accepted $\gamma\gamma \rightarrow \tau^+\tau^-$ events is given as a function of the absolute value of the transverse momentum sum of the two charged particles and the tagged electron, normalised to the beam energy: $|\sum \vec{p}_t|/E_{\text{beam}}$. The points represent the data, the open histogram the Vermaseren Monte Carlo normalised to the integrated luminosity. The shaded histogram represents the same quantity for the single-tag $\gamma\gamma \rightarrow \mu^+\mu^-$ events, with an arbitrary normalisation. The striking difference between the $\gamma\gamma \rightarrow \mu^+\mu^-$ and $\gamma\gamma \rightarrow \tau^+\tau^-$ events

can be explained by the presence of neutrinos and neutral pions in the decays of the tau leptons, which are absent in the $\gamma\gamma \rightarrow \mu^+\mu^-$ events. This difference and the fact that the $\gamma\gamma \rightarrow \tau^+\tau^-$ distribution is well reproduced by the Monte Carlo confirm that the data are consistent with a significant production rate of $\gamma\gamma \rightarrow \tau^+\tau^-$ events. In Figure 6b we show the distribution of the accepted $\gamma\gamma \rightarrow \tau^+\tau^-$ events as a function of the Q^2 of the photon, tagged by the electron observed in the forward detector. The measured values (data points) are in reasonable agreement with the Monte Carlo expectation (solid line).

9 Conclusions

We have observed the reaction $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$, both in the single-tag and in the double-tag mode. We have extracted the photon structure function F_2 from the single-tag events and found agreement with the QED expectation. We have selected 1462 single-tag events with an estimated background of 24.0 ± 2.4 events, and 48 double-tag events with an estimated background of 1.7 ± 0.6 events, which should be compared with the Monte Carlo expectations of 1418 ± 142 single-tag and 36.9 ± 5.0 double-tag events, respectively. We have identified 34.9 ± 6.7 events of the reaction $\gamma\gamma \rightarrow \tau^+\tau^-$ in the single-tag mode on an estimated background of 5.1 ± 2.7 events. This is in good agreement with a QED Monte Carlo simulation using the Vermaseren generator, which predicts 32.7 ± 4.8 events. The Q^2 distribution of the $\gamma\gamma \rightarrow \tau^+\tau^-$ events agrees with the expectation from the Vermaseren Monte Carlo.

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